

Example 5

Propagation of a Beam of Light in the radial-direction of Sirius from the Gravitational Radius of Sirius until the Earth

with the speed of light $c = \frac{1}{\sqrt{\epsilon_0 \mu_0}}$

within a Radial Gravitational Field $G1$ oriented in the radial-direction of Sirius with acceleration $G1 = g[r]$, resulting in Gravitational Redshift $\nu_{\text{Grav}-1}$ (GRS) caused by the: Gravitational Field of Sirius.

There are **two** kinds of Redshift caused by the gravitational field of the sun.

- 1) Redshift₁ caused by the Gravitational Field within the **fusion core** of the sun (this is the area where the emitted light of the sun has been created). This has been calculated in Example 3
- 2) Redshift₂ caused by the Gravitational Field outside the fusion core of the sun. This will be calculated in Example 4

The total RedShift caused by the Gravitational Field of the sun is the sum of Redshift₁ + Redshift₂

Book : ***Rising of the James Webb Space Telescope and its Fundamental Blindnes*** : Page 46,

Equation (73)

Example of a Spherical Beam of Light , propagating in the radial – direction (r – direction)

with the speed of light $c = \frac{1}{\sqrt{\epsilon_0 \mu_0}}$.

The input for the Electric Field Intensity $\{E(x, y, z) = ev\}$ and the Magnetic Field Intensity $\{H(x, y, z) = mv\}$ have been

substituted in the Field Equation for the Electromagnetic Field within a constant gravitational field with acceleration g in the z – direction. (**Book Equation 73, page 46**).

$$-\frac{1}{c^2} \frac{\partial(\mathbf{E} \times \mathbf{H})}{\partial t} + \epsilon_0 \mathbf{E} (\nabla \cdot \mathbf{E}) - \epsilon_0 \mathbf{E} \times (\nabla \times \mathbf{E}) + \mu_0 \mathbf{H} (\nabla \cdot \mathbf{H}) - \mu_0 \mathbf{H} \times (\nabla \times \mathbf{H}) + \frac{1}{2} \left(\epsilon^2 \mu (\mathbf{E} \cdot \mathbf{E}) + \epsilon \mu^2 (\mathbf{H} \cdot \mathbf{H}) \right) \mathbf{g} = 0 \quad \text{Equation (73)}$$

Which can be written as :

$$\text{term1} = -\frac{1}{c^2} \frac{\partial(\mathbf{E} \times \mathbf{H})}{\partial t}$$

$$\text{term2} = \epsilon_0 \mathbf{E} (\nabla \cdot \mathbf{E})$$

$$\text{term3} = -\epsilon_0 \mathbf{E} \times (\nabla \times \mathbf{E})$$

$$\text{term4} = \mu_0 \mathbf{H} (\nabla \cdot \mathbf{H})$$

$$\text{term5} = -\mu_0 \mathbf{H} \times (\nabla \times \mathbf{H})$$

$$\text{term6} = -\frac{1}{2} \left(\epsilon^2 \mu (\mathbf{E} \cdot \mathbf{E}) + \epsilon \mu^2 (\mathbf{H} \cdot \mathbf{H}) \right) \mathbf{g}$$

Equation (73) equals :

$$\text{term1} + \text{term2} + \text{term3} + \text{term4} + \text{term5} + \text{term6} = 0$$

$$\ln[*]= \epsilon_0 =.$$

$$\ln[*]= \mu_0 =.$$

$$\ln[*]= \mathbf{x} =.$$

$$\ln[*]= \mathbf{y} =.$$

In[*]:= $Z = .$

In[*]:= $t = .$

In[*]:= Get["VectorAnalysis`"]

⋯ General: VectorAnalysis` is now obsolete. The legacy version being loaded may conflict with current functionality. See the Compatibility Guide for updating information.

In[*]:= Get["Calculus`DSolve`"]

⋯ Get: Cannot open Calculus`DSolve`.

Out[*]:= **\$Failed**

In[*]:= InverseFunctions → True

Out[*]:= InverseFunctions → True

In[*]:= Needs["DifferentialEquations`NDSolveProblems`"]

In[*]:= Needs["DifferentialEquations`NDSolveUtilities`"]

In[*]:= Needs["DifferentialEquations`InterpolatingFunctionAnatomy`"]

In[*]:= Get["Calculus`DSolveIntegrals`"]

⋯ Get: Cannot open Calculus`DSolveIntegrals`.

Out[*]:= **\$Failed**

In[*]:= SetCoordinates[Cartesian[x, y, z]]

Out[*]:= Cartesian [x, y, z]

In[*]:= $G1 = .$

In[*]:= $\epsilon 0 = .$

In[*]:= $\mu 0 = .$

In[*]:= $fg = .$

In[*]:= $M_{\text{sun}} = .$

The light being emitted by the sun has been created by **nuclear fusion in the core of the the sun**.

Outside the **Gravitational Radius** of the Sun, (the distance of the center of the sun where the gravitational acceleration of the sun changes from increasing into decreasing) , the generated light will propagate with the speed of light :

$$c = \frac{1}{\sqrt{\epsilon_0 \mu_0}}$$

oppositely directed towards the gravitational acceleration outside the sun. The gravitational acceleration outside the sun will increase proportional to :

$$g^2[r] = K^2 \frac{1}{r^2} = \frac{f_g M_{\text{Sun}}}{4\pi} \frac{1}{r^2}$$

within a **Radial Gravitational Field** $g[r]$ oriented in the radial –

direction of the **Sun** with **acceleration** $g[r]$ resulting in the Gravitational Redshift caused by the **Gravitational Field outside the Sun**.

The **Gravitational Intensity Shift** for the light emitted by the sun equals according the Book: **Page 64, Equation (111)**

$$I_{\text{NGR}} = e^{-\frac{K^2 \epsilon_0 \mu_0}{z}}$$

Instead of Spherical Coordinates (r, θ, φ) for the light emitted Spherically within the sun, for the far field (within a small area like the dimensions of a **Space Telescope**), Cartesian Coordinates (x, y, z) are being used.

$$\text{In[*]}:= \epsilon 0 = .$$

$$\text{In[*]}:= \mu 0 = .$$

$$\text{In[*]}:= fg = .$$

$$\text{In[*]}:= \text{Msun} = .$$

$$\text{In[*]}:= K1 = .$$

$$\text{In[*]}:= K2 = .$$

$$\text{In[*]}:= h[z] = K3 e^{-\frac{K2 \epsilon 0 \mu 0}{2z}}$$

$$\text{Out[*]}:= e^{-\frac{K2 \epsilon 0 \mu 0}{2z}} K3$$

$$\text{In[*]}:= \text{ev} = \{h[z] \times g[(t - z \sqrt{\epsilon 0} \sqrt{\mu 0})], 0, 0\}$$

$$\text{Out[*]}:= \left\{ e^{-\frac{K2 \epsilon 0 \mu 0}{2z}} K3 g[t - z \sqrt{\epsilon 0} \sqrt{\mu 0}], 0, 0 \right\}$$

$$\text{In[*]}:= \text{mv} = (1/\text{Sqrt}[\mu 0]) * \text{Sqrt}[\epsilon 0] * \{0, h[z] \times g[(t - z \sqrt{\epsilon 0} \sqrt{\mu 0})], 0\}$$

$$\text{Out[*]}:= \left\{ 0, \frac{e^{-\frac{K2 \epsilon 0 \mu 0}{2z}} K3 \sqrt{\epsilon 0} g[t - z \sqrt{\epsilon 0} \sqrt{\mu 0}]}{\sqrt{\mu 0}}, 0 \right\}$$

$$\text{In[*]}:= g2 = \left\{ 0, 0, -\frac{K2}{z^2} \right\}$$

$$\text{Out[*]}:= \left\{ 0, 0, -\frac{K2}{z^2} \right\}$$

$$\text{In[*]}:= \text{Intensity} = \frac{1}{2} (\epsilon 0 (\text{Dot}[\text{ev}, \text{ev}]) + \mu 0 (\text{Dot}[\text{mv}, \text{mv}]))$$

$$\text{Out[*]}:= \frac{1}{2} \left(e^{-\frac{K2 \epsilon 0 \mu 0}{z}} K3^2 \epsilon 0 g[t - z \sqrt{\epsilon 0} \sqrt{\mu 0}]^2 + e^{-\frac{K2 \epsilon 0 \mu 0}{z}} K3^2 g[t - z \sqrt{\epsilon 0} \sqrt{\mu 0}]^2 \epsilon 0 \right)$$

$$\text{In[*]}:= \text{FullSimplify}[\%]$$

$$\text{Out[*]}:= \frac{e^{-\frac{K2 \epsilon 0 \mu 0}{z}} K3^2 g[t - z \sqrt{\epsilon 0} \sqrt{\mu 0}]^2 1 (\epsilon 0 + \epsilon 0)}{2}$$

$$\text{In[*]}:= \text{Div}[\text{ev}]$$

$$\text{Out[*]}:= 0$$

$$\text{In[*]}:= \text{Div}[\text{mv}]$$

$$\text{Out[*]}:= 0$$

$$\text{In[*]}:= \text{FullSimplify}[\%]$$

$$\text{Out[*]}:= 0$$

$$\text{In[*]:= term1a} = \text{D}[\text{Cross}[\text{ev}, \text{mv}], \text{t}]$$

$$\text{Out[*]:= } \left\{ \theta, \theta, \frac{2 e^{-\frac{K2 \epsilon_0 \mu_0}{z}} K3^2 \sqrt{\epsilon_0} \mathbf{g}[\mathbf{t} - z \sqrt{\epsilon_0} \sqrt{\mu_0}] \mathbf{g}'[\mathbf{t} - z \sqrt{\epsilon_0} \sqrt{\mu_0}]}{\sqrt{\mu_0}} \right\}$$

$$\text{term1} = -\frac{1}{c^2} \frac{\partial(\mathbf{E} \times \mathbf{H})}{\partial t}$$

$$\text{In[*]:= term1} = ((-\epsilon_0) * \mu_0) * \text{D}[\text{Cross}[\text{ev}, \text{mv}], \text{t}]$$

$$\text{Out[*]:= } \left\{ \theta, \theta, -2 e^{-\frac{K2 \epsilon_0 \mu_0}{z}} K3^2 \epsilon_0^{3/2} \sqrt{\mu_0} \mathbf{g}[\mathbf{t} - z \sqrt{\epsilon_0} \sqrt{\mu_0}] \mathbf{g}'[\mathbf{t} - z \sqrt{\epsilon_0} \sqrt{\mu_0}] \right\}$$

$$\text{term2} = \epsilon_0 \mathbf{E} (\nabla \cdot \mathbf{E})$$

$$\text{In[*]:= term2} = \epsilon_0 * \text{ev} * \text{Div}[\text{ev}]$$

$$\text{Out[*]:= } \{ \theta, \theta, \theta \}$$

$$\text{term3} = -\epsilon_0 \mathbf{E} \times (\nabla \times \mathbf{E})$$

$$\text{In[*]:= term3} = (-\epsilon_0) * \text{Cross}[\text{ev}, \text{Curl}[\text{ev}]]$$

$$\text{Out[*]:= } \left\{ \theta, \theta, -\epsilon_0 \left(\frac{e^{-\frac{K2 \epsilon_0 \mu_0}{z}} K2 K3^2 \epsilon_0 \mu_0 \mathbf{g}[\mathbf{t} - z \sqrt{\epsilon_0} \sqrt{\mu_0}]^2}{2 z^2} - e^{-\frac{K2 \epsilon_0 \mu_0}{z}} K3^2 \sqrt{\epsilon_0} \sqrt{\mu_0} \mathbf{g}[\mathbf{t} - z \sqrt{\epsilon_0} \sqrt{\mu_0}] \mathbf{g}'[\mathbf{t} - z \sqrt{\epsilon_0} \sqrt{\mu_0}] \right) \right\}$$

$$\text{term4} = \mu_0 \mathbf{H} (\nabla \cdot \mathbf{H})$$

$$\text{In[*]:= term4} = \mu_0 * \text{mv} * \text{Div}[\text{mv}]$$

$$\text{Out[*]:= } \{ \theta, \theta, \theta \}$$

$$\text{term5} = -\mu_0 \mathbf{H} \times (\nabla \times \mathbf{H})$$

$$\text{In[*]:= term5} = (-\mu_0) * \text{Cross}[\text{mv}, \text{Curl}[\text{mv}]]$$

$$\text{Out[*]:= } \left\{ \theta, \theta, -\mu_0 \left(\frac{e^{-\frac{K2 \epsilon_0 \mu_0}{z}} K2 K3^2 \epsilon_0^2 \mathbf{g}[\mathbf{t} - z \sqrt{\epsilon_0} \sqrt{\mu_0}]^2}{2 z^2} - \frac{e^{-\frac{K2 \epsilon_0 \mu_0}{z}} K3^2 \epsilon_0^{3/2} \mathbf{g}[\mathbf{t} - z \sqrt{\epsilon_0} \sqrt{\mu_0}] \mathbf{g}'[\mathbf{t} - z \sqrt{\epsilon_0} \sqrt{\mu_0}]}{\sqrt{\mu_0}} \right) \right\}$$

$$\text{term6} = -\frac{1}{2} \left(\epsilon^2 \mu \left(\overline{E \cdot E} \right) + \epsilon \mu^2 \left(\overline{H \cdot H} \right) \right) \overline{g}$$

$$\text{In[*]:= term6} = -\left(\frac{\epsilon^2 \mu}{2} \text{Dot}[ev, ev] + \frac{\epsilon \mu^2}{2} \text{Dot}[mv, mv] \right) \overline{g}$$

$$\text{Out[*]:= } \left\{ 0, 0, \frac{e^{-\frac{K2 \epsilon \mu}{z}} K2 K3^2 \epsilon^2 \mu \overline{g} [t - z \sqrt{\epsilon} \sqrt{\mu}]^2}{z^2} \right\}$$

In[*]:= vergelijking = term1 + term2 + term3 + term4 + term5 + term6

$$\begin{aligned} \text{Out[*]:= } & \left\{ 0, 0, \frac{e^{-\frac{K2 \epsilon \mu}{z}} K2 K3^2 \epsilon^2 \mu \overline{g} [t - z \sqrt{\epsilon} \sqrt{\mu}]^2}{z^2} - \right. \\ & 2 e^{-\frac{K2 \epsilon \mu}{z}} K3^2 \epsilon^{3/2} \sqrt{\mu} \overline{g} [t - z \sqrt{\epsilon} \sqrt{\mu}] \overline{g}' [t - z \sqrt{\epsilon} \sqrt{\mu}] - \\ & \mu \left(\frac{e^{-\frac{K2 \epsilon \mu}{z}} K2 K3^2 \epsilon^2 \overline{g} [t - z \sqrt{\epsilon} \sqrt{\mu}]^2}{2 z^2} - \right. \\ & \left. \left. \frac{e^{-\frac{K2 \epsilon \mu}{z}} K3^2 \epsilon^{3/2} \overline{g} [t - z \sqrt{\epsilon} \sqrt{\mu}] \overline{g}' [t - z \sqrt{\epsilon} \sqrt{\mu}]}{\sqrt{\mu}} \right) - \right. \\ & \left. \epsilon \left(\frac{e^{-\frac{K2 \epsilon \mu}{z}} K2 K3^2 \epsilon \mu \overline{g} [t - z \sqrt{\epsilon} \sqrt{\mu}]^2}{2 z^2} - \right. \right. \\ & \left. \left. e^{-\frac{K2 \epsilon \mu}{z}} K3^2 \sqrt{\epsilon} \sqrt{\mu} \overline{g} [t - z \sqrt{\epsilon} \sqrt{\mu}] \overline{g}' [t - z \sqrt{\epsilon} \sqrt{\mu}] \right) \right\} \end{aligned}$$

In[*]:=

The electromagnetic force density in the x - direction equals :

$$\text{In[*]:= xv} \text{vergelijking} = \text{term1}[[1]] + \text{term2}[[1]] + \text{term3}[[1]] + \text{term4}[[1]] + \text{term5}[[1]] + \text{term6}[[1]]$$

$$\text{Out[*]:= } 0$$

$$\text{In[*]:= FullSimplify}[\%]$$

$$\text{Out[*]:= } 0$$

$$\text{In[*]:= xv} \text{vergelijking1} = \%$$

$$\text{Out[*]:= } 0$$

In[*]:=

The electromagnetic force density in the y - direction equals :

$$\text{In[*]:= y} \text{vergelijking} = \text{term1}[[2]] + \text{term2}[[2]] + \text{term3}[[2]] + \text{term4}[[2]] + \text{term5}[[2]] + \text{term6}[[2]]$$

$$\text{Out[*]:= } 0$$

In[*]:= FullSimplify[%]

Out[*]= 0

In[*]:= yvergelijking1 = %

Out[*]= 0

The electromagnetic force density in the z - direction equals :

In[*]:= zvergelijking = term1[[3]] + term2[[3]] + term3[[3]] + term4[[3]] +
term5[[3]] + term6[[3]]

$$\text{Out[*]} = \frac{e^{-\frac{K_2 \epsilon_0 \mu_0}{z}} K_2 K_3^2 \epsilon_0^2 \mu_0 g[t - z \sqrt{\epsilon_0} \sqrt{\mu_0}]^2}{z^2} -$$

$$2 e^{-\frac{K_2 \epsilon_0 \mu_0}{z}} K_3^2 \epsilon_0^{3/2} \sqrt{\mu_0} g[t - z \sqrt{\epsilon_0} \sqrt{\mu_0}] g'[t - z \sqrt{\epsilon_0} \sqrt{\mu_0}] -$$

$$\mu_0 \left(\frac{e^{-\frac{K_2 \epsilon_0 \mu_0}{z}} K_2 K_3^2 \epsilon_0^2 g[t - z \sqrt{\epsilon_0} \sqrt{\mu_0}]^2}{2 z^2} - \right.$$

$$\left. \frac{e^{-\frac{K_2 \epsilon_0 \mu_0}{z}} K_3^2 \epsilon_0^{3/2} g[t - z \sqrt{\epsilon_0} \sqrt{\mu_0}] g'[t - z \sqrt{\epsilon_0} \sqrt{\mu_0}]}{\sqrt{\mu_0}} \right) -$$

$$\epsilon_0 \left(\frac{e^{-\frac{K_2 \epsilon_0 \mu_0}{z}} K_2 K_3^2 \epsilon_0 \mu_0 g[t - z \sqrt{\epsilon_0} \sqrt{\mu_0}]^2}{2 z^2} - \right.$$

$$\left. e^{-\frac{K_2 \epsilon_0 \mu_0}{z}} K_3^2 \sqrt{\epsilon_0} \sqrt{\mu_0} g[t - z \sqrt{\epsilon_0} \sqrt{\mu_0}] g'[t - z \sqrt{\epsilon_0} \sqrt{\mu_0}] \right)$$

In[*]:= FullSimplify[%]

Out[*]= 0

In[*]:= zvergelijking1 = %

Out[*]= 0

Results for the electromagnetic force densities in resp x-direction, y-direction, z-direction:

In[*]:= xvergelijking1

Out[*]= 0

In[*]:= yvergelijking1

Out[*]= 0

In[*]:=

Out[*]= 0

According the force-density equations in the x-direction, y-direction and z-direction, the resulting electromagnetic force density equals zero in every direction. This represents the solution for equation

(73) on page 47.

It follows from the mathematical solution for the electromagnetic field:

$$\mathbf{e}_v = e^{-\frac{K_2 \epsilon_0 \mu_0}{z}} \mathbf{g} \left[e^{-\frac{K_2 \epsilon_0 \mu_0}{2z}} (t - z \sqrt{\epsilon_0} \sqrt{\mu_0}) \right]$$

that the **intensity increases** with the value:

$$\text{Intensity} = \frac{1}{2} (\epsilon_0 (\text{Dot}[\mathbf{e}_v, \mathbf{e}_v]) + \mu_0 (\text{Dot}[\mathbf{m}_v, \mathbf{m}_v])) =$$

$$= e^{-\frac{K_2 \epsilon_0 \mu_0}{z}} \epsilon_0 \mathbf{g} \left[e^{-\frac{K_2 \epsilon_0 \mu_0}{2z}} (t - z \sqrt{\epsilon_0} \sqrt{\mu_0}) \right]^2$$

The Electromagnetic Energy Intensity is proportional to: $e^{-\frac{K_2 \epsilon_0 \mu_0}{z}}$ along the distance z in the direction **opposite** to the z -direction of the gravitational field.

The frequency is proportional to the energy density (**book: Equation 98 Page 55**) and the wavelength is inversely proportional to the energy. The speed of light remains **constant** in a gravitational field.

Within a Gravitational Field $\frac{K_2}{z^2}$ The observed Cosmological Redshift will be:

$$\omega_{\text{NLGR}} = \omega_0 e^{-\frac{K_2 \epsilon_0 \mu_0}{z}} \sqrt{\frac{1 + \frac{v}{c}}{1 - \frac{v}{c}}}$$

$$\text{The term } \left(\frac{\Delta \omega}{\omega} = e^{-\frac{K_2 \epsilon_0 \mu_0}{z}} \right)$$

has been presented in generally as a **GRS** Redshift comparable with the Doppler Shift generated by a velocity v_{Doppler} :

$$v_{\text{Doppler}} = c \frac{\Delta \omega}{\omega} [\text{m s}^{-1}] = c e^{-\frac{K2 \epsilon_0 \mu_0}{z}}$$

There are two types of **GRS** generated by the sun.

1) GRS generated inside the sun where the gravitational field is proportional to the radial distance z

$g[z] = K1 z$ with the Solution:

$$v_{\text{Doppler-1}} = c \frac{\Delta \omega}{\omega} [\text{m s}^{-1}] = c e^{-\frac{1}{4} K1 z^2 \mu_0 \epsilon_0}$$

(book, page 64, Equation (111))

2) GRS generated outside the sun where the gravitational field is proportional to the radial distance $\frac{1}{z^2}$

$g[z] = K2 \frac{1}{z^2}$ with the Solution:

$$v_{\text{Doppler-2}} = c \frac{\Delta \omega}{\omega} [\text{m s}^{-1}] = c e^{-\frac{K2 \epsilon_0 \mu_0}{z}}$$

(book, page 64, Equation (111))

$$\text{In[*]:= Intensity} = -\frac{1}{2} (\epsilon_0 (\text{Dot}[ev, ev]) + \mu_0 (\text{Dot}[mv, mv]))$$

$$\text{Out[*]:=} -\frac{1}{2} 1 \left(e^{-\frac{K2 \epsilon_0 \mu_0}{z}} K3^2 \epsilon_0 g[t - z \sqrt{\epsilon_0} \sqrt{\mu_0}]^2 + e^{-\frac{K2 \epsilon_0 \mu_0}{z}} K3^2 g[t - z \sqrt{\epsilon_0} \sqrt{\mu_0}]^2 \epsilon_0 \right)$$

In[*]:= FullSimplify[%]

$$\text{Out[*]:=} -\frac{e^{-\frac{K2 \epsilon_0 \mu_0}{z}} K3^2 g[t - z \sqrt{\epsilon_0} \sqrt{\mu_0}]^2 1 (\epsilon_0 + \epsilon_0)}{2}$$

2) GRS generated outside the sun where the gravitational field is proportional

to the radial distance $\frac{1}{z^2}$

$g[z] = K2 \frac{1}{z^2}$ with the Solution:

$$v_{\text{Doppler-2}} = c \frac{\Delta \omega}{\omega} [\text{m s}^{-1}] = c e^{-\frac{K2 \epsilon_0 \mu_0}{z}}$$

(book, page 64, Equation (111))

$$v_{\text{Doppler-1}} = c \frac{\Delta \omega}{\omega} [\text{m s}^{-1}] = c e^{-\frac{K2 \epsilon_0 \mu_0}{z2}} - c e^{-\frac{K2 \epsilon_0 \mu_0}{z1}}$$

With:

$z_1 =$ Radius of the sun

$z_2 =$ Distance between the sun and the Earth

$$\text{In[*]}:= c = 3 \times 10^8$$

$$\text{Out[*]}:= 300\,000\,000$$

$$\text{In[*]}:= \epsilon_0 = 8.85 \times 10^{-12}$$

$$\text{Out[*]}:= 8.85 \times 10^{-12}$$

$$\text{In[*]}:= \mu_0 = 4 \pi 10^{-7}$$

$$\text{Out[*]}:= \frac{\pi}{2500000}$$

$$\text{In[*]}:= fg = 6.67428 \times 10^{-11}$$

$$\text{Out[*]}:= 6.67428 \times 10^{-11}$$

$$\text{In[*]}:= \text{Msun} = 1.98892 \times 10^{30}$$

$$\text{Out[*]}:= 1.98892 \times 10^{30}$$

$$\text{In[*]}:= \text{MSiriusB} = 1.018 \text{ Msun}$$

$$\text{Out[*]}:= 2.02472 \times 10^{30}$$

$$\text{In[*]}:= R_{\text{Sun}} = 69\,634\,000$$

$$\text{Out[*]}:= 69\,634\,000$$

$$\text{In[*]}:= R_{\text{SiriusB}} = 1.711 R_{\text{Sun}}$$

$$\text{Out[*]}:= 1.19144 \times 10^8$$

$$\text{In[*]}:= \text{Lightyear} = 9.461 \times 10^{15}$$

$$\text{Out[*]}:= 9.461 \times 10^{15}$$

$$\text{In[*]}:= z2 = 8.611 \text{ Lightyear}$$

$$\text{Out[*]}:= 8.14687 \times 10^{16}$$

$$\text{In[*]}:= K2 = \frac{fg \text{ MSiriusB}}{4 \pi}$$

$$\text{Out[*]}:= 1.07537 \times 10^{19}$$

$$\text{In[*]}:= z1 = 0.0063875 R_{\text{Sun}}$$

$$\text{Out[*]}:= 444\,787.$$

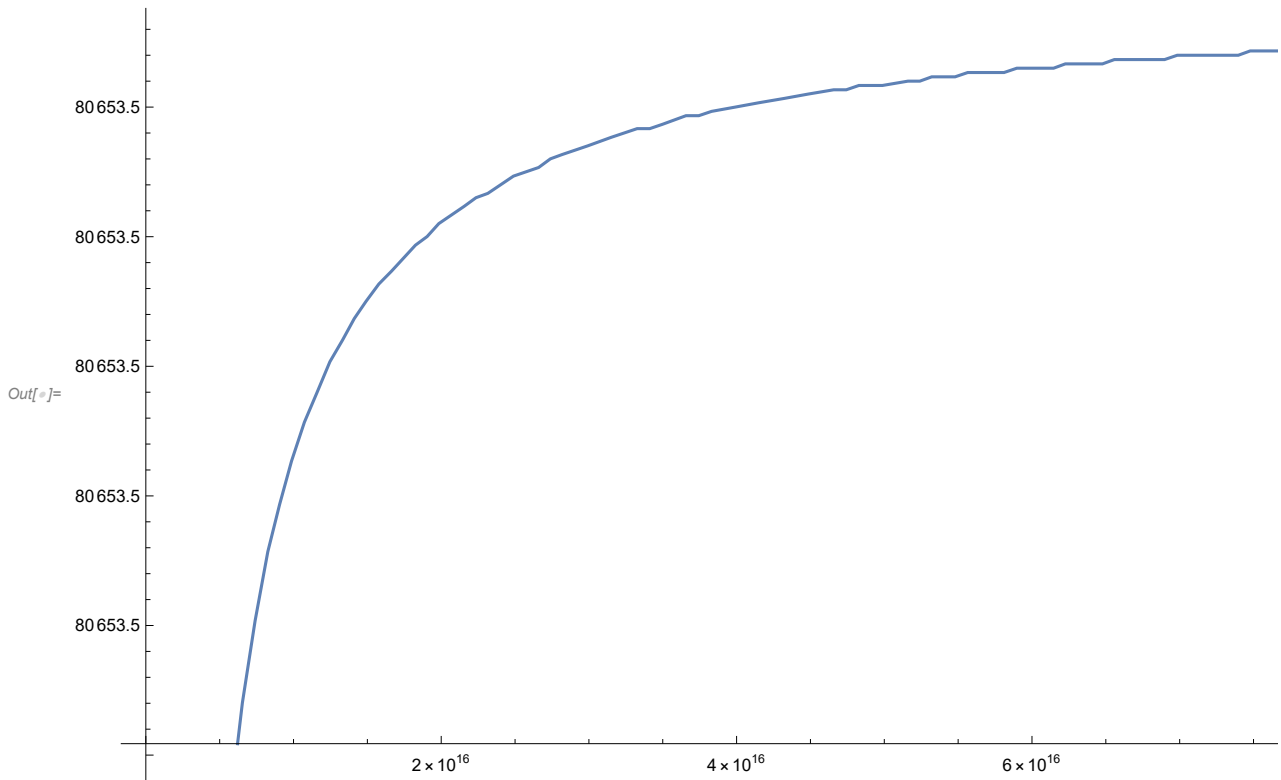
The distance between Sirius and Earth equals z2

In[*]:= z2 = 8.611 Lightyear

Out[*]:= 8.14687 × 10¹⁶

The Gravitational RedShift from the “Intense Emitting Radiation Radius z1” of Sirius has been presented below :

In[*]:= Plot[c (e^{- $\frac{K2 \epsilon_0 \mu_0}{z}$} - e^{- $\frac{K2 \epsilon_0 \mu_0}{z1}$}), {z, z1, z2}]



In[*]:= VDoppler = c (e^{- $\frac{K2 \epsilon_0 \mu_0}{z2}$} - e^{- $\frac{K2 \epsilon_0 \mu_0}{z1}$})

Out[*]:= 80 653.5

The calculated value of 80.653 k[m/s] for the Gravitational Redshift for Sirius corresponds

to the measured average value for the **GRS** for Sirius of = 80.65 + /-0.77 km/s.
published in :

The gravitational redshift of Sirius B